

# **Lunar perturbation of the metric associated to the averaged orbital transfer**

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# Riemannian metric for the unperturbed problem

Controlled Kepler equation  
(planar case)

$$\frac{dx}{dt} = \sum_{i=1}^2 u_i F_i(x, l)$$

$$\frac{dl}{dt} = w_0(x, l) + g(x, l, u).$$

$x = (n, \rho, \theta)$  are the slow orbital elements of the satellite.

$n$  is the mean motion,

$\rho$  the eccentricity,

$\theta$  the argument of periapsis,

$l$  is the fast angle pointing out the satellite's position on his orbit.

Averaged Hamiltonian for the energy cost

$$\bar{H} = \frac{1}{4n^{\frac{5}{3}}} \left[ 18n^2 p_n^2 + 5(1 - \rho^2) p_\rho^2 + (5 - 4\rho^2) \frac{p_\theta^2}{\rho^2} \right].$$

The Riemannian metric associated

$$g = \frac{1}{9n^{\frac{1}{3}}} dn^2 + \frac{2n^{\frac{5}{3}}}{5(1 - \rho^2)} d\rho^2 + \frac{2n^{\frac{5}{3}}}{5 - 4\rho^2} d\theta^2.$$

# Model for the lunar perturbation

The perturbed dynamic can be written as

$$\frac{dx}{dt} = F_0(x, l, l') + \sum_{i=1}^2 u_i F_i(x, l)$$

$$\frac{dl}{dt} = \tilde{w}_0(x, l) + g(x, l, u)$$

$$\frac{dl'}{dt} = w_1(l')$$

adding  $l'$  as a fast angle pointing out the satellite's position on his orbit.

The double averaged drift vector field corresponding to the lunar perturbation is

$$\overline{F}_0(n, p, \theta) = \frac{3n'^2}{4n} \sqrt{1 - \rho^2} \frac{\partial}{\partial \theta} (n, p, \theta).$$

# Zermelo navigation problem

A Zermelo navigation problem on a  $n$ -dimensional Riemannian manifold  $(\mathcal{X}, g)$  is a **time minimal** problem associated to the system

$$\frac{dx}{dt} = X_0(x) + \sum_{i=1}^n u_i X_i(x)$$

where  $F_i$  form an orthonormal frame for the metric  $g$  and  $|u| \leq 1$ .

The double averaged Zermelo Hamiltonian associated to our controlled system is

$$\begin{aligned} H_{perb} &= \langle p, \bar{F}_0(x, p) \rangle + \varepsilon \sqrt{\bar{H}(x, p)} \\ &= p_\theta \left[ \frac{3n'^2}{4n} \sqrt{1 - \rho^2} \right] + \varepsilon \sqrt{\frac{1}{4n^{\frac{5}{3}}} \left[ 18n^2 p_n^2 + 5(1 - \rho^2) p_\rho^2 + (5 - 4\rho^2) \frac{p_\theta^2}{\rho^2} \right]}. \end{aligned}$$

# Geometric concept of conjugate point

Let  $\vec{H}_{pert}$  the Hamiltonian vector field associated to  $H_{pert}$ .

-  $z = (x, p)$  is a reference extremal, solution of  $\vec{H}_{pert}$  on  $[0, t_f]$ .

Jacobi equation

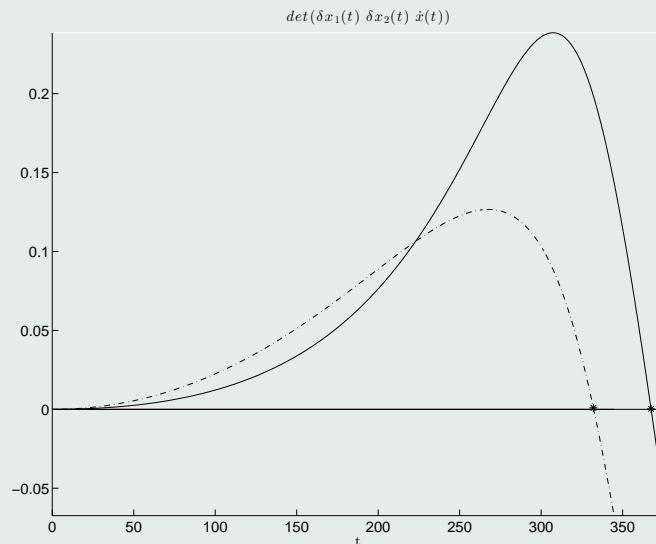
$$\dot{\delta z}(t) = d\vec{H}_{pert}(z(t))\delta z(t)$$

A Jacobi field is a non trivial solution  $\delta z = (\delta x, \delta p)$ , vertical at time  $t$  if  $\delta x(t) = 0$ .

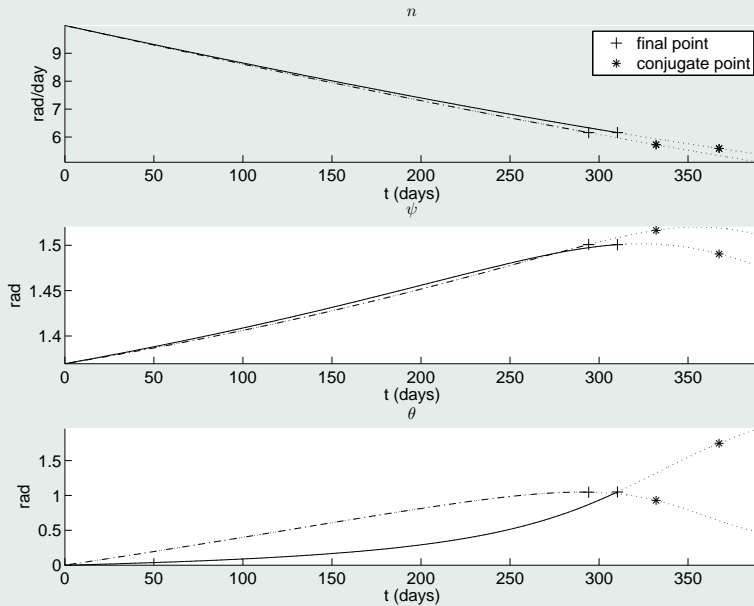
-  $t_c$  : the first conjugate time such that the map  $p_0 \rightarrow \exp_{x_0}(t, p_0) = x(t, x_0, p_0)$  is not an immersion at  $t = t_c$ .

- Rank test condition to compute  $t_c$  for a given trajectory  $x(t, x_0, p_0)$

$$\det[\delta x_1(t_c), \dots, \delta x_{n-1}(t_c), \dot{x}(t_c)] = 0.$$

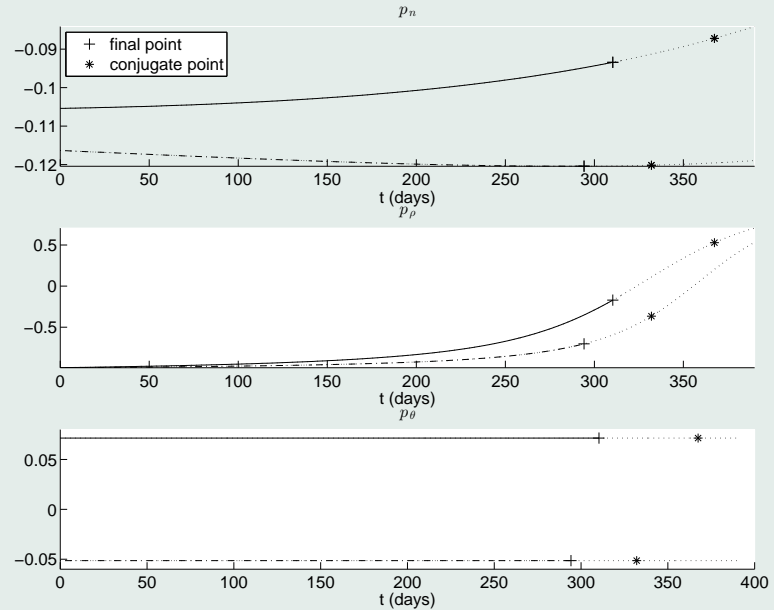


# Time evolution of trajectories



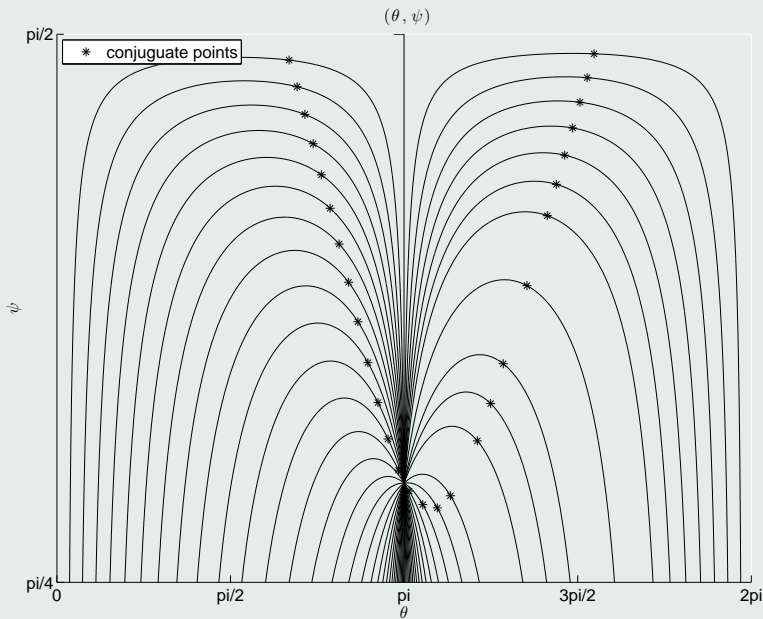
*Evolution of the state vector*  
 $(n(t), \psi(t) = \frac{\pi}{2} - \text{asin}(\rho(t)), \theta(t)).$

solide line : unperturbed case,  
dash-dot line : perturbed case.

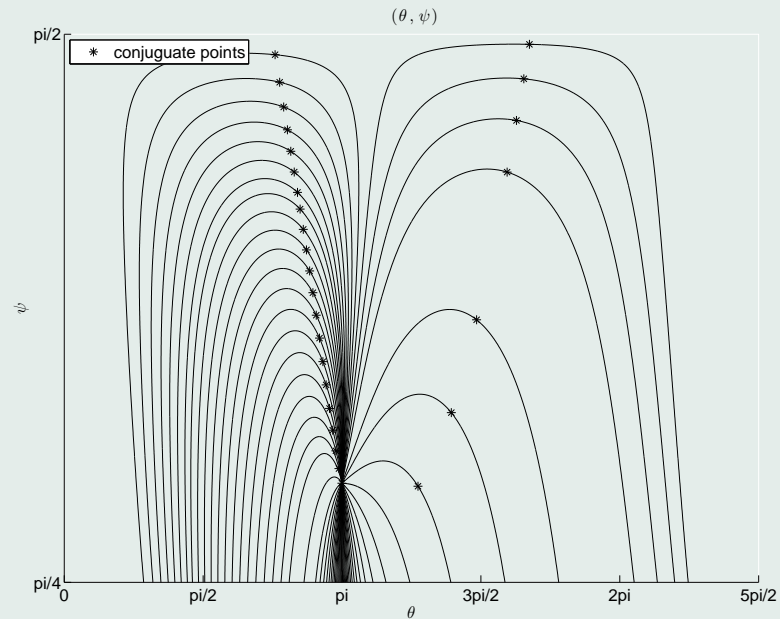


*Evolution of the adjoint vector*  
 $(p_n(t), p_\rho(t), p_\theta(t)).$

# Trajectories in $(\theta, \psi)$ coordinates



*Trajectories from the unperturbed case.*



*Trajectories from the perturbed case.*